Weighted Evolving Networks

S. H. Yook, H. Jeong, and A.-L. Barabási

Department of Physics, University of Notre Dame, Notre Dame, Indiana 46556

Y. Tu

IBM, TJ Watson Research Center, POB 218, Yorktown Heights, New York 10598 (Received 9 January 2001)

Many biological, ecological, and economic systems are best described by weighted networks, as the nodes interact with each other with varying strength. However, most evolving network models studied so far are binary, the link strength being either 0 or 1. In this paper we introduce and investigate the scaling properties of a class of models which assign weights to the links as the network evolves. The combined numerical and analytical approach indicates that asymptotically the total weight distribution converges to the scaling behavior of the connectivity distribution, but this convergence is hampered by strong logarithmic corrections.

DOI: 10.1103/PhysRevLett.86.5835

PACS numbers: 89.75.Hc, 02.50.Cw, 84.35.+i, 87.23.Ge

The recognition that real networks are fundamentally different from the random models that dominated the mathematical literature in the past 40 years [1,2] leads to a surge of activity in addressing the statistical properties of these systems [3-10]. In one aspect most recently developed models, aimed to describe the large-scale topology of complex networks, are incomplete when compared with real systems: they assume that all links are equivalent. But in many fields it is well known that the interaction strengths can vary widely, such variations being essential to the network's ability to carry on its basic functions. Sociologists have repeatedly argued about the importance of assigning strengths to social links, finding that the weak links people have outside their close circle of friends play a key role in keeping the social system together [11]. Recently, Newman has showed that assigning weights to the links between scientists allows for a better characterization of the scientific collaboration web [12]. Similarly, there is an ongoing discussion about the importance of weak links between species in guaranteeing the stability of an ecosystem [13]. Finally, many transportation networks in nature, ranging from cardiovascular to respiratory networks, have well defined weights or flow rates assigned to the links, whose magnitude is intimately determined by the network's topology [14]. The issue of link strength has been extensively addressed in the neural network literature. The question posed in that context so far had a unique focus: given a network topology, how can one alter the link weights in a dynamical fashion to allow the network to perform certain desired functions, ranging from memory to pattern recognition [15]? Similarly, research on allometric scaling has also been concerned with assigning weights to links on a network with fixed, often treelike topology [14]. On the other hand, the recent advances in statistical modeling of complex networks have brought the community's attention towards large networks whose topology evolves in time. Despite the known importance of interaction strengths in the various systems these models aim to describe, in this context there have been no attempts to model networks other than binary nets, whose links have weights 0 or 1.

In this paper we take a first step in the direction of a systematic study of evolving networks with nonbinary connectivities. We introduce and investigate two models that assign weights to new links as they are dynamically created, providing a prototype of a weighted evolving network. While we choose the simplest possible models, in which the weights are driven by the network connectivity only, numerical simulations indicate that the distribution of the total weight scales differently from the total connectivity. However, an analytical solution reveals that the different scaling behavior can be explained by strong logarithmic correction, and asymptotically the investigated weighted networks belong to the same universality class as their unweighted counterparts.

Weighted scale-free (WSF) model.—Starting from a small number (m_0) of vertices, at each time step we add a new node which links to *m* existing nodes in the system. The probability that a new node *j* will connect to an existing node *i* is

$$\Pi_i = \frac{k_i}{\sum_j k_j},\tag{1}$$

where k_i is the total number of links that the node *i* has. In assigning a weight to the newly established link $j \leftrightarrow i$, we assume that the weight w_{ji} (= w_{ij}) is proportional to k_i , i.e., more connected (and therefore more "powerful") nodes gain more weight. Also, one can assume that all new nodes have fairly uniform total "resources" for linking to other nodes in the system. We therefore require that each new node has a fixed total weight; i.e., we normalize w_{ij} such that the sum of the weights for the *m* new links is $\sum_{\{i'\}} w_{ji'} = 1$, where $\{i'\}$ represents a sum over the *m* existing nodes to which the new node *j* is connected. As a result of the two assumptions, each link $i \leftrightarrow j$ of the newly added node *j* is assigned a weight as

$$w_{ji} = \frac{k_i}{\sum_{\{i'\}} k_{i'}}.$$
 (2)

Weighted exponential (WE) model.—The model is inspired by model A discussed in Refs. [16,17], and is defined as follows: at every time step we add a new node with $m (\leq m_0)$ links, connected with equal probability to the nodes present in the system. The weights of the links are assigned again by using (2).

The difference between the WSF and WE models comes in preferential attachment, which is known to fundamentally alter the topology [7–9,16–18]: The WSF model generates a scale-free network whose connectivity distribution follows $P(k) \sim k^{-3}$, while the network generated by the WE model is exponential with the connectivity distribution following $P(k) = \frac{e}{m}e^{-k/m}$. Since the weights of the links are driven by the connectivity, this difference is expected to lead to significant changes in the distribution of the link strengths as well.

We start by investigating the weight distribution of the two models. As Figs. 1(a) and 1(b) show, both the WE and the WSF models lead to a peaked and skewed weight distribution, whose tails decay exponentially (or faster) for large w_{ij} . The boundedness of $P(w_{ij})$ is due to the normalization condition, which does not allow individual weights to be larger than 1. Most important, however, we find that the distribution is stationary; i.e., $P(w_{ij})$ is independent of time (and system size).

While the individual weights assigned to links, w_{ij} , are bounded, we get a very different picture when we study the total weight associated with a selected node. In binary networks the node's importance is characterized by the total number of links it has, k_i . Similarly, in a weighted network the importance of a node *i* can be measured by its total weight, obtained by summing the weights of the links that connect to it, $w_i = \sum_{\{j\}} w_{ij}$.



FIG. 1. The distribution $P(w_{ij})$ of the individual link weights, w_{ij} , for the (a) WE and the (b) WSF models, defined in the text (m = 2). The symbols correspond to different system sizes (or time), i.e., $N = 10^3$ (\bigcirc), 10^4 (\square), 10^5 (\diamondsuit), and 10^6 (\triangle). The insets shows the same data on a log-linear plot, indicating that the tail decays faster than exponential.

Because of the normalization condition (2) a new node has $w_i = 1$, but w_i increases in time every time when subsequently added nodes link to i. Since in both models the weights are determined by the network connectivity, we expect that P(w) closely follows P(k). In contrast, the numerical results summarized in Fig. 2 indicate striking differences between P(k) and P(w). As Fig. 2(a) shows, while for the WE model P(k) decays exponentially, P(w)systematically deviates from a simple exponential behavior. This difference is even more evident in the network dynamics: while both $k_i(t)$ and $w_i(t)$ appear to increase logarithmically in time, they can be fitted with a different slope on a log-linear plot [Fig. 2(b)]. Similar systematic discrepancies are observed for the WSF model as well: as Fig. 2c indicates, while P(w) and P(k) can be fitted with power laws, $P(w) \sim w^{-\sigma}$ and $P(k) \sim k^{-\gamma}$, it appears that $\gamma = 3$ and the exponent σ is different from γ . Furthermore, we find that σ depends strongly on m [Fig. 2(c)]. Again, this difference is reflected in the dynamical behavior of $k_i(t)$ and $w_i(t)$: as Fig. 2(d) indicates, $w_i(t) \sim t^{\beta}$ with $\beta > 1/2$, in contrast with $k_i(t) \sim t^{1/2}$ [6,17] predicted by the binary scale-free model.

To understand the different behaviors of w_i and k_i uncovered by the numerical simulations, we resort to the analytical method in determining the averaged behavior of $w_i(t)$ for the discussed model. To simplify the discussion in the following we assume m = 2; however, the calculations can be generalized for arbitrary m. The total weight of node i at time t can be written as



FIG. 2. (a) Distribution P(w) of the total connectivity w assigned to individual nodes for the WE model. The symbols correspond to different values of m, i.e., m = 2 (\bigcirc), 3 (\square), 4 (\diamondsuit), and 5 (\triangle). The inset shows the connectivity distribution, P(k), for the same parameters as in the main panel. (b) Time dependence of $k_i(t)$ (\bigcirc) and $w_i(t)$ (\square) for a randomly selected node i for the WE model (i = 5000). (c) P(k) (\bigcirc) and P(w) (\square) distributions for the WSF model for m = 5. The inset shows the same data for m = 2. (d) $k_i(t)$ (\bigcirc), $w_i(t)$ (\square) vs t for the WSF model (i = 10000).

$$w_i(t) = 1 + \sum_{\{j\}} w_{ij} = 1 + \int_{t_i^0}^t \tilde{P}_i(t') \langle w_{ij}(t') \rangle dt', \quad (3)$$

where $\tilde{P}_i(t)$ is the probability that node *i* is selected to be connected to a new node *j* at time *t* and t_i^0 is the time at which the node *i* has been added to the system. $\langle w_{ij} \rangle$ is the average weight of link $i \leftrightarrow j$ once the link is established. When a new node *j* and the list of *m* nodes $\{i'\}$ to which it connects are selected, the weights of the links, $w_{ji'}$, are assigned according to (2). These weights depend on the number of links the selected nodes have, i.e., $\{k_{i'}\}$. If we assume that node *j* is connected to nodes *i* and *l* (*m* = 2), we have

$$\langle w_{ij}(t) \rangle = \int_m^\infty w_{ji}(l) \mathcal{P}(k_l) \, dk_l \,,$$
 (4)

where $w_{ji}(l)$ is the weight between the *j* and *i* nodes, $\mathcal{P}(k_l)$ is the probability distribution of k_l , the total link number of node *l*. Substituting (4) into (3), we obtain

$$w_{i}(t) = 1 + \int_{t_{i}^{0}}^{t} \int_{m}^{\infty} \tilde{P}_{i}(t') w_{ji}(l) \mathcal{P}(k_{l}) dk_{l} dt'.$$
 (5)

According to (2) for m = 2, the weight $w_{ji}(l)$ is given by

$$w_{ji}(l) = \frac{k_i}{k_i + k_l},\tag{6}$$

thus Eq. (5) becomes

$$w_i(t) = 1 + \int_{t_i^0}^t \int_m^\infty \tilde{P}_i(t') \frac{k_i}{k_i + k_l} \,\mathcal{P}(k_l) \,dk_l \,dt'.$$
(7)

Equation (7) represents a general expression for calculating $w_i(t)$ for m = 2. To apply it to the WE and WSF models, we need to calculate explicitly $\tilde{P}(t)$ and $\mathcal{P}(k_l)$.

WE model.—In the WE model the nodes to which a new node connects are selected uniformly among all existing nodes, thus the probability that node i will be picked is independent of this node's connectivity and is given by

$$\tilde{P}_i(t) = \frac{m}{t+m_0}.$$
(8)

Similarly, the connectivity distribution and the dynamical behavior of a single node are given by [17]

$$\mathcal{P}(k) = Ae^{-k/m} = \frac{e}{m}e^{-k/m},$$

$$k_i(t) = m[\ln(m_0 + t - 1) - \ln(m_0 + t_i^0 - 1) + 1]$$

$$= m[\ln(at + b) + 1],$$
(9)

where $a = 1/(m_0 + t_i^0 - 1)$, $b = (m_0 - 1)/(m_0 + t_i^0 - 1)$, and the normalization condition is $1 = \int_m^\infty \mathcal{P}(k) dk$.

Substituting (9) into (7), we obtain

$$w_i(t) = 1 + e \int_{t_i^0}^t \int_m^\infty \frac{1}{t' + m_0} \frac{k_i(t')}{k_i(t') + k_l} e^{-k_l/m} dk_l dt'.$$

After performing the integration and inserting $k_i(t)$ from (9), for large t we obtain

$$w_i(t) \simeq m \ln(at + b) - m \ln[\ln(at + b) + 2] + C,$$
(10)

where *C* is an integration constant independent of *t*. Therefore the relation between $w_i(t)$ and $k_i(t)$ for large *t* follows:

$$w_i(t) \simeq k_i(t) - m \ln \ln t + C.$$
⁽¹¹⁾

The prediction (11) is fully supported by numerical simulations: in Fig. 3(a) we plot the difference $w_i(t) - k_i(t)$ as a function of $\ln \ln(t)$, showing that the difference indeed follows a double logarithmic law. This result is very interesting since it indicates that the different slopes observed in Fig. 2(b) for $k_i(t)$ and $w_i(t)$ do not represent distinct power law scaling behaviors, but are the result of logarithmic corrections.

WSF model.—In the scale-free model the probability distributions and $k_i(t)$ are given by [17]

$$\tilde{P}_{i}(t) = m \frac{k_{i}(t)}{\sum_{j}^{t} k_{j}} = m \frac{k_{i}(t)}{2mt} = \frac{k_{i}(t)}{2t},$$

$$\mathcal{P}(k) = mk^{-2} \quad [\alpha \ k \cdot P(k)],$$

$$k_{i}(t) = \frac{m}{\sqrt{t_{i}^{0}}} \sqrt{t}.$$
(12)

Substituting (12) into (7), and performing the integrals we obtain

$$w_i(t) \simeq k_i(t) - \frac{m}{8} \left(\ln \frac{m^2 t}{t_i^0} \right)^2 + \frac{m}{2} \ln(m) \ln\left(\frac{t}{t_i^0}\right) + C',$$
(13)

indicating that despite a different scaling behavior suggested by the numerical simulations [Fig. 2(d)], we are dealing with strong logarithmic corrections and asymptotically two scaling laws are the same. Again, the analytical prediction (13) is confirmed by more detailed numerical simulations shown in Fig. 3(b).

Our ability to calculate analytically w_{ij} for the discussed models is based on the fact that the weights are driven by



FIG. 3. The difference $[w_i(t) - k_i(t)]$ for the (a) WE and the (b) WSF models. The continuous lines in each case represent the analytic solutions (11) and (13), respectively. We limited the simulations to nodes appearing at large t_i^0 ($t_i^0 = 10^4$) to capture the asymptotic limit that is predicted by our predictions (11) and (13). We find that for smaller t_i^0 the crossover time for the convergence to the analytic solution is numerically prohibited.

the connectivity distribution. To address the generality of our results we investigated several extensions of these two models that we discuss in the following.

Weight driven weight.—In general, one could expect that in some systems the quantity determining the weight is not the connectivity, but are the weights themselves. To investigate this possibility we replaced (2) with

$$w_{ji} = \frac{w_i}{\sum_{i'} w_{i'}};$$
 (14)

i.e., the weight of the newly added links are determined by the total weight of the nodes. While we cannot solve this model analytically, the numerical results are similar to those observed for the WE and WSF models: an apparently different scaling behavior for k and w can be attributed to slow corrections to scaling.

Weight driven connectivity.—In some systems the topology could be driven by the total weights, and not by the connectivity. Thus we assume that the probability (1) that a new node is connected to a node j is

$$\Pi_i = \frac{w_i}{\sum_j w_j},\tag{15}$$

where w_i is the weight of node *i*. The weights are then assigned following (2). We find that the scaling of this network is identical to that of the scale-free model, and the evolution of the weights also follows the paradigm established for the WSF model.

Discussion.—Extensive simulations of networks whose size is comparable to the real networks that are currently available indicate the emergence of new scaling exponents for the behavior of the total weights. However, the analytical solutions reveal that the results are affected by strong logarithmic corrections, and asymptotically the scaling behaviors of the weighted and unweighted models are identical. This result raises important questions regarding our ability to uncover the correct scaling behavior of real weighted networks, should such data become available in the near future: the real exponents could be easily shadowed by corrections to scaling similar to that encountered in the investigated models here.

The results presented in this paper represent only the starting point towards understanding weighted networks. In some real systems, diverse dynamical rules can govern the assignment of weights to links, which could result in statistical properties of the network that are different from that discussed here. In particular, we assumed that once a weight has been assigned to a link, it stays unchanged, which is often not the case in more realistic networks: weights can evolve dynamically just as the network topology does. For example, acquaintance can turn into friendship by strengthening a previously weak link. Determining the generic behavior of such complex evolving systems is a real challenge for future research. Despite these limitations, the investigated models give a glimpse into the complex behavior we are facing as we attempt to make network modeling more realistic by incorporating weights.

Research at Notre Dame supported by NSF, PHY-9988674, and CAREER DMR97-01998.

- [1] B. Bollobás, Random Graphs (Academic, London, 1985).
- [2] P. Erdös and A. Rényi, Publ. Math. Inst. Hung. Acad. Sci. 5, 17 (1960).
- [3] D. J. Watts, Small Worlds: The Dynamics of Networks Between Order and Randomness (Princeton University Press, Princeton, 1999).
- [4] D. J. Watts and D. H. Strogatz, Nature (London) 393, 440 (1998).
- [5] M.E.J. Newman, J. Stat. Phys. 101, 819 (2000).
- [6] A.-L. Barabási and R. Albert, Science 286, 509 (1999).
- [7] S.N. Dorogovtsev and J.F.F. Mendes, Phys. Rev. E 62, 1842 (2000); Europhys. Lett. 52, 33 (2000).
- [8] P. L. Krapivsky, S. Redner, and F. Leyvraz, Phys. Rev. Lett. 85, 4629 (2000).
- [9] L. A. N. Amaral, A. Scalar, M. Barthélémy, and H. E. Stanley, Proc. Natl. Acad. Sci. U.S.A. 97, 11 149 (2000).
- [10] R. Cohen and K. Erez, D. ben-Avraham, and S. Havlin, Phys. Rev. Lett. 85, 4626 (2000).
- [11] M. Granovetter, Am. J. Soc. 78, 1360 (1973).
- [12] M.E.J. Newman, Proc. Natl. Acad. Sci. U.S.A. 98, 404 (2001).
- [13] E.L. Berlow, Nature (London) 398, 330 (1999).
- [14] J.R. Banavar, A. Maritan, and A. Rinaldo, Nature (London) 399, 130 (1999).
- [15] B. Müller and J. Reinhardt, Neural Networks: An Introduction (Springer-Verlag, Berlin, 1991).
- [16] R. Albert, H. Jeong, and A.-L. Barabási, Nature (London) 401, 130 (1999).
- [17] A.-L. Barabási, R. Albert, and H. Jeong, Physica (Amsterdam) 272A, 173 (1999).
- [18] L. Kullmann and J. Kertész, cond-mat/0012410.